

Spreadsheet Approach to Partial Differential Equations. Part 2: Parabolic and Hyperbolic Equations*

C. Y. LAM

School of Mechanical and Production Engineering, Nanyang Technological University, Nanyang Avenue, Singapore 2263

As a companion paper to Part 1 [1], three other interactive spreadsheet programs developed by using the Lotus 1-2-3 spreadsheet software to solve the parabolic one-dimensional diffusion equation and the hyperbolic one-dimensional wave equation are presented. For the diffusion equation, the explicit forward time central space method and the implicit Crank-Nicholson method have been implemented in the first two programs. For the wave equation, the explicit central difference method has been implemented in the third program. The advanced interactive graphics capability is demonstrated in two of these programs employing the explicit methods. It is found that this feature is very useful in performing the 'What-if' analysis graphically. These programs have been tested in the classroom and favourable comments have been obtained.

INTRODUCTION

THE solution of the elliptic Laplace equation by using the spreadsheet approach has been presented in Part 1 [1]. Apart from the elliptic equation, the parabolic and hyperbolic equations are also the core discussion topics in engineering and science subjects related to partial differential equations. Generally, the parabolic one-dimensional diffusion equation and the hyperbolic one-dimensional wave equation are used to illustrate the numerical methods in the teaching of these subjects. The reasons are that these equations are simple and that they have a wide variety of physical applications; for example, the heat conduction (diffusion), the nuclear diffusion and the unsteady fluid flow problems for the diffusion equation and the vibration, the electricity transmission and the supersonic flow problems for the wave equation. The finite difference methods that are usually included in teaching are the explicit forward time central space (FTCS) method and the Crank-Nicolson (CN) method for the diffusion equation and the explicit central difference (CD) method for the wave equation. These methods are based on a marching technique in the time direction. Here, three programs developed to implement these finite difference methods by using the spreadsheet approach are described. The objective is again to provide the students with better alternative tools (other than the traditional computer programs), which can be used to enhance their understanding of the numerical methods concerned. The usefulness of the advanced interactive graphics feature

provided by the Lotus 1-2-3 [2], which is impossible to achieve by the traditional approach, is explored and is found to be extremely helpful in viewing the effects of varying the input parameters in graphical form automatically for the explicit methods.

PROBLEM FORMULATION

The finite difference methods used in the spreadsheet programs can be found in many standard texts, for example [3-5].

The one-dimensional diffusion equation

The parabolic one-dimensional diffusion equation is

$$\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2} \quad (1)$$

where c is a constant and c^2 is the diffusivity, u is the dependent variable, t and x are the time and spatial coordinates respectively.

To solve the diffusion equation by finite difference methods, the physical t - x domain is discretized and the grid points are labelled as shown in Fig. 1. To obtain the numerical solution, two Dirichlet boundary conditions ($u(0, t)$ and $u(L, t), t \geq 0$) and one initial condition ($u(x, 0), 0 < x < L$) are specified. The solution procedure is based on the marching technique in the t -direction.

The FTCS method. This method is also known as the explicit Euler method. By using the forward difference and the central difference representations for the time and spatial derivatives, respect-

* Paper accepted 3 August 1992.

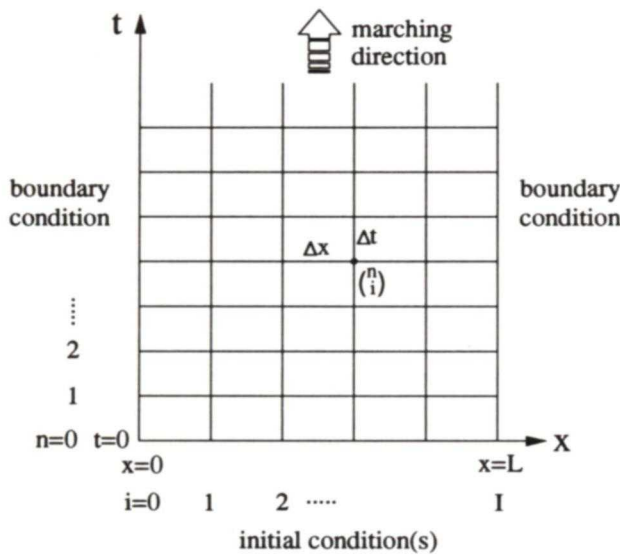


Fig. 1. The discretized computational domain

ively, the finite difference equation for equation (1) is

$$u_i^{n+1} = F(u_{i-1}^n + u_{i+1}^n) + (1 - 2F)u_i^n \quad (2)$$

where $F = c^2\Delta t/(\Delta x)^2$, known as the grid Fourier number, and the subscripts and the superscripts denote the numbers of the x - and t -grid points respectively. In this explicit method, $F \leq 1/2$ for stability.

The CN method. In this method, the spatial derivative is replaced by the mean of the central differences at the two consecutive t -grid points (n and $n + 1$) and the time derivative is replaced by the central difference in the middle between these two grid points. The finite difference equation for equation (1) is

$$-Fu_{i-1}^{n+1} + (2 + 2F)u_i^{n+1} - Fu_{i+1}^{n+1} = Fu_{i-1}^n + (2 - 2F)u_i^n + Fu_{i+1}^n \quad (3)$$

where all parameters are the same as those defined for equation (2).

Applying this at the interior grid points $i = 1, \dots, I - 1$ for marching one t -step from n to $n + 1$ yields

$$\begin{bmatrix} 2+2F & -F & 0 & 0 & 0 & \dots & 0 & 0 \\ -F & 2+2F & -F & 0 & 0 & \dots & 0 & 0 \\ 0 & -F & 2+2F & -F & 0 & \dots & 0 & 0 \\ 0 & 0 & -F & 2+2F & -F & \dots & 0 & 0 \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ 0 & 0 & 0 & 0 & 0 & \dots & -F & 2+2F \end{bmatrix} \begin{bmatrix} u_1^{n+1} \\ u_2^{n+1} \\ u_3^{n+1} \\ u_4^{n+1} \\ \cdot \\ \cdot \\ u_{I-1}^{n+1} \end{bmatrix} = \begin{bmatrix} Fu_0^n + (2 - 2F)u_1^n + Fu_2^n + Fu_0^{n+1} \\ Fu_1^n + (2 - 2F)u_2^n + Fu_3^n \\ Fu_2^n + (2 - 2F)u_3^n + Fu_4^n \\ Fu_3^n + (2 - 2F)u_4^n + Fu_5^n \\ \cdot \\ \cdot \\ Fu_{I-2}^n + (2 - 2F)u_{I-1}^n + Fu_I^n + Fu_{I-1}^{n+1} \end{bmatrix} \quad (4)$$

where I is the number of x -intervals.

This implicit method is unconditionally stable.

The one-dimensional wave equation

The hyperbolic one-dimensional wave equation is

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} \quad (5)$$

where c (= constant) is the wave propagation speed, the other parameters are the same as those defined for equation (1).

In applying the finite difference method, the $x-t$ domain is discretized and the grid points are labelled in the same manner as for the diffusion equation shown in Fig. 1. However, to obtain the numerical solution, two Dirichlet boundary conditions (u_0^n and $u_I^n, n \geq 0$) and two initial conditions (u_i^0 for $0 < i < I$ and $\partial u_i^0/\partial t$ for $0 \leq i \leq I$) are specified. Similar to that of the diffusion equation, the solution procedure is based on marching in the t -direction.

The CD method. Using the central differences for both derivatives, the finite difference equation for the wave equation is

$$u_i^{n+1} = -u_i^{n-1} + R^2(u_{i+1}^n + u_{i-1}^n) + 2(1 - R^2)u_i^n \quad (6a)$$

where $R = c\Delta t/\Delta x$, known as the Courant number, and the subscripts and superscripts denote the numbers of the x - and t -grid points respectively.

Equation (6a) is used for $n > 0$. For the first time step, $n = 0$ and u_i^{-1} in equation (6a) is replaced by that obtained from the initial condition $\partial u_i^0/\partial t$ using the central difference and the equation becomes

$$u_i^1 = \frac{R^2}{2} (u_{i+1}^0 + u_{i-1}^0) + (1 - R^2)u_i^0 + (\Delta t) \frac{\partial u_i^0}{\partial t} \quad (6b)$$

where $\partial u_i^0/\partial t = \partial u(x_i, 0)/\partial t$ is prescribed.

This method is explicit and $R \leq 1$ for stability.

THE SPREADSHEET PROGRAMS

The programs developed are menu driven with a common menu almost identical to that for the programs for solving Laplace equation as presented in Part 1 (see fig. 3 and table 1 in [1]). The only differences are that the Graph T-variation command replaces the Graph Y-variation command because t is now one of the independent variables and that the Data Use command will continue the calculations at more t -steps (rather than the iteration) using the existing data since the finite difference methods used here are based on the time marching procedure. The same user-friendly and interactive features as those described in Part 1 [1]

are retained. The only difference is in the procedure of using these programs due to the time marching procedure. In this case, after the marching of a t -step is completed, the user will be requested to indicate whether the marching is to be continued for the next t -step. If the response is negative, the program will return to the main menu.

In the programs, the step sizes Δx and Δt are computed automatically based on the input data L , I , F and R .

The interactive graphics feature. The Lotus 1-2-3 Release 3.1 offers a useful interactive graphics capability which is ideal in studying the effects on the solution graphically when an input parameter is changed—the so-called ‘What-if’ analysis. This is done by using the two simple 1-2-3 commands /Worksheet Window Graph and /Worksheet Window Clear. First, the program command Quit must be used to return to 1-2-3’s READY mode after a solution has been obtained and a graph has been created by the program command Graph X-variation or Graph T-variation. To set the interactive graphics mode, the arrow keys are used to move the 1-2-3 cursor to a column which can be any one apart from the leftmost column shown on the screen. Upon invoking the 1-2-3’s /Worksheet Window Graph command, the monitor screen will be split vertically at that column and the current graph will be displayed on the right side of the screen. Some of the input data can then be altered manually (by moving the 1-2-3 cursor to the respective location(s) followed by keying in the new value(s) and hitting the Enter key) and the numerical results as well as the current graph will be updated automatically. The interactive graphics mode can be cleared by using the 1-2-3’s /Worksheet Window Clear command.

The interactive graphics mode as explained can be used only when the spreadsheet cells storing the solution contain formulae which depend algebrai-

cally on the contents of the other cells storing the input data so that the contents of the former cells as well as the graph will be updated accordingly when the contents of the latter cells are changed. For the programs employing the FTCS method for the diffusion equation and the CD method for the wave equations, any input data except the number of x -intervals I can be altered manually after a solution has been obtained. The value of I cannot be altered in this manner because new x -grid points are required for a different value of I and the program can only generate new x -grid points within the command Data Use. For the CN method implemented for the diffusion equation, the interactive graphics mode cannot be used since each t -step marching is done by solving equation (4) using matrix operations in a way that the contents of the cells storing the solution are numerical values.

All programs assume constant Dirichlet boundary conditions. The boundary conditions are required to be entered only once at the starting step corresponding to $t = 0$ and they will be used for the time marching in subsequent t -steps. However, for the FTCS method and the CD method, variable Dirichlet boundary conditions can also be used as will be discussed later. The initial condition is entered at each interior grid point corresponding to $t = 0$ accordingly.

The spreadsheet program FTCS.WK3

This spreadsheet program employs the FTCS method to solve the diffusion equation. To march one t -step, the program assigns equation (2) at the new grid points and the numerical results are obtained. This is done by the 1-2-3’s /Copy command programmed in the macros.

Example. Consider the unsteady one-dimensional heat conduction problem of a thin metal rod perfectly insulated laterally of length $L = 1$ m, $c = 1.1$ m/s^{0.5} and initial temperature

```

A:H20: 10
Any changes for the boundary or initial conditions? <Y/N> : _
READY
A      A      B      C      D      E      F      G      H
1
2 =====
3 SOLVING THE 1-D DIFFUSION EQUATION BY THE FTCS METHOD
4 Conditions : constant steps in x and t, constant B.C.'s (may not equal)
5 [© C Y Lam, Nanyang Technological University]
6 =====
7 Equation to be solved :
8 du(x,t)/dt = c^2 * d^2u(x,t)/dx^2 where d denotes partial differentiation
9
10          Constant c, c =          1.1
11          Solution domain in x, L =          1
12          Number of x intervals, I =          6
13          Grid Fourier number, F =          0.3
14          Step in x, dx =          0.166667
15          Step in t, dt =          0.006887
16
17 RESULTS
18          x=          x=          x=          x=          x=          x=
19          t          0 0.166667 0.333333          0.5 0.666667 0.833333          1
20 -----
FTCS.WK3
CMD
    
```

Fig. 2. The input data in FTCS.WK3

$u(x, 0) = 0^\circ\text{C}$. At $t = 0$ s, the temperature at one end $u(0, t)$ is suddenly brought to 30°C and the temperature at the other end $u(1, t)$ is brought to 10°C .

Using the program command Data Input, the data are entered as shown in Fig. 2 in which $I = 6$ and $F = 0.3$. In this figure, the x -grid points have been overlaid on the working sheet A automatically before the boundary and initial conditions are entered along the first row corresponding to $t = 0$. The results for 15 t -steps are obtained as shown in Fig. 3(a) and the variations of u with x are graphed as shown in Fig. 3(b).

Figure 4 shows the screen in interactive graphics mode invoked by the 1-2-3's command /Worksheet Window Graph with the cursor moved to

column D. It can be seen that the screen is divided vertically with the current graph (variations with x in this case) displayed on the right. To study the effects of varying the grid Fourier number F , the content of the cell with address A:E12 which contains F is changed manually and the numerical results as well as the graph are updated automatically as shown in Fig. 5(a-d) for $F = 0.5, 0.51, 0.55$ and 0.6 respectively. These figures clearly verify that the FTCS method is unstable for $F > 0.5$.

The 'What-if' analysis is also done by varying the other parameters. Figure 6 shows the results by changing the boundary condition $u(0, t)$ to 10°C and Fig. 7 shows the results by changing the initial condition $u(x, 0)$ to $-20 \sin(\pi x)$. In Fig. 7, the new initial condition is entered as 1-2-3 formulae

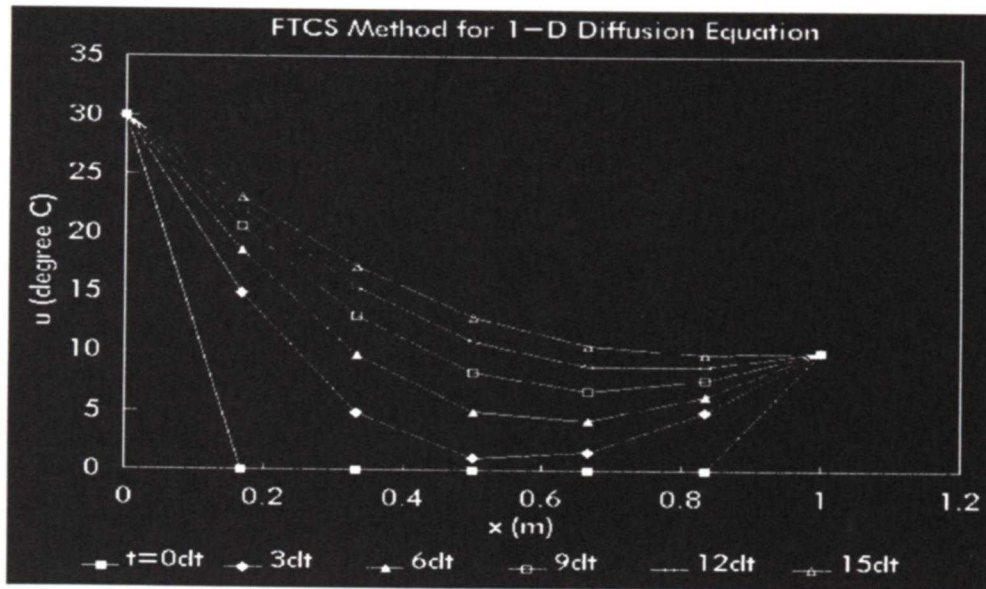
A:A35: +A34+\$DT
Proceed another t-step ? <Y/N> : _

READY

| A | B | C | D | E | F | G | H | |
|----|----------|----|----------|----------|----------|----------|----------|---|
| 16 | RESULTS | | | | | | | |
| 17 | | x= | x= | x= | x= | x= | x= | |
| 18 | t | 0 | 0.166667 | 0.333333 | 0.5 | 0.666667 | 0.833333 | 1 |
| 19 | | | | | | | | |
| 20 | 0 | 30 | 0 | 0 | 0 | 0 | 10 | |
| 21 | 0.006887 | 30 | 9 | 0 | 0 | 0 | 10 | |
| 22 | 0.013774 | 30 | 12.6 | 2.7 | 0 | 0.9 | 4.2 | |
| 23 | 0.020661 | 30 | 14.85 | 4.86 | 1.08 | 1.62 | 4.95 | |
| 24 | 0.027548 | 30 | 16.398 | 6.723 | 2.376 | 2.457 | 5.466 | |
| 25 | 0.034435 | 30 | 17.5761 | 8.3214 | 3.7044 | 3.3354 | 5.9235 | |
| 26 | 0.041322 | 30 | 18.52686 | 9.71271 | 4.9788 | 4.22253 | 6.37002 | |
| 27 | 0.048209 | 30 | 19.32456 | 10.93678 | 6.172092 | 5.093658 | 6.814767 | |
| 28 | 0.055096 | 30 | 20.01086 | 12.02371 | 7.277969 | 5.933521 | 7.254004 | |
| 29 | 0.061983 | 30 | 20.61146 | 12.99613 | 8.298356 | 6.733 | 7.681658 | |
| 30 | 0.068871 | 30 | 21.14342 | 13.8714 | 9.238082 | 7.487204 | 8.092563 | |
| 31 | 0.075758 | 30 | 21.61879 | 14.66301 | 10.10281 | 8.194075 | 8.483187 | |
| 32 | 0.082645 | 30 | 22.04642 | 15.38168 | 10.89825 | 8.85343 | 8.851497 | |
| 33 | 0.089532 | 30 | 22.43307 | 16.03607 | 11.62983 | 9.466296 | 9.196628 | |
| 34 | 0.096419 | 30 | 22.78405 | 16.6333 | 12.30264 | 10.03446 | 9.51854 | |
| 35 | 0.103306 | 30 | 23.10361 | 17.17933 | 12.92139 | 10.56014 | 9.817753 | |

FTCS.WK3

Fig. 3. The solution of the one-dimensional heat conduction problem for $I = 6$ and $F = 0.3$ using FTCS.WK3
(a) The solution for 15 t -steps



(b) Variations of u with x as displayed on the screen

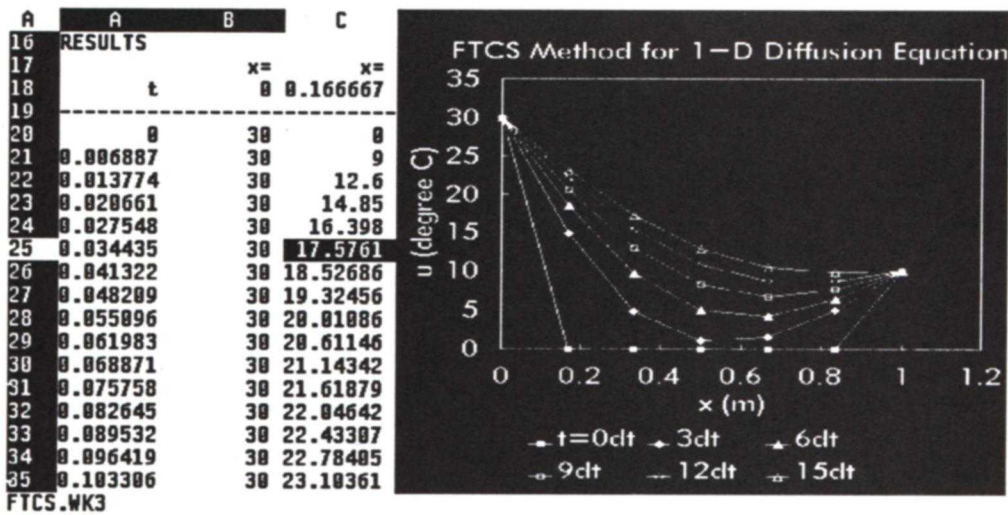


Fig. 4. The interactive graphics screen for $I = 6$ and $F = 0.3$ using FTCS.WK3

A:E12: 0.5

READY

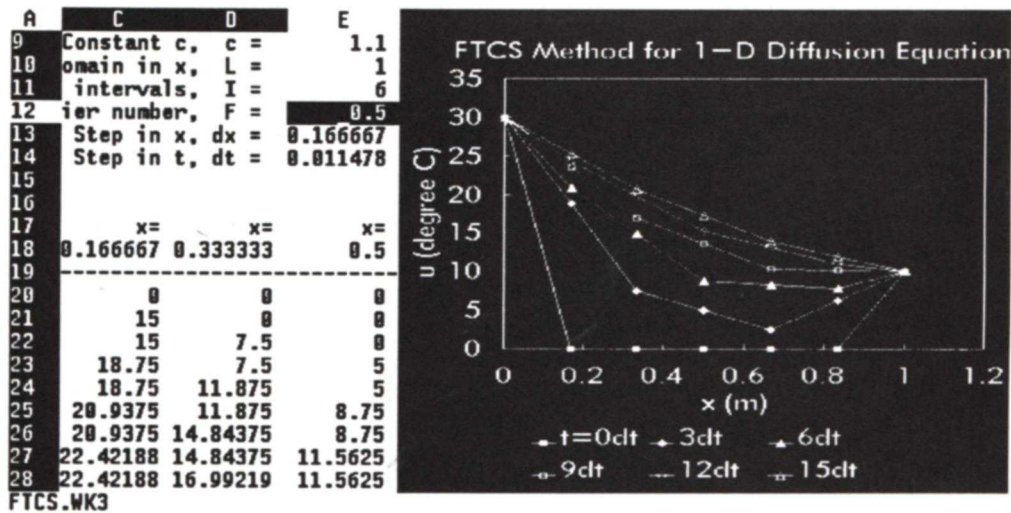
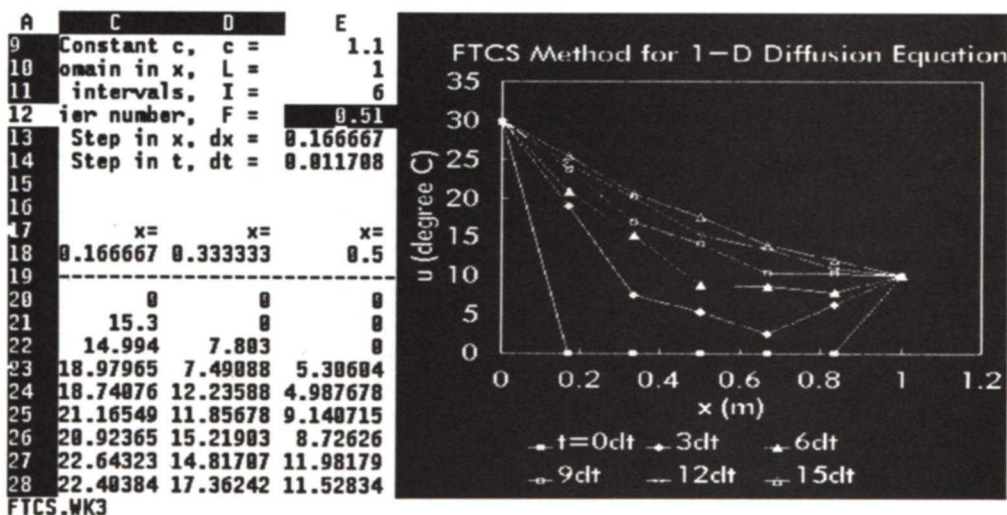


Fig. 5. The effects of varying F using FTCS.WK3 ($I = 6$)
(a) $F = 0.5$

A:E12: 0.51

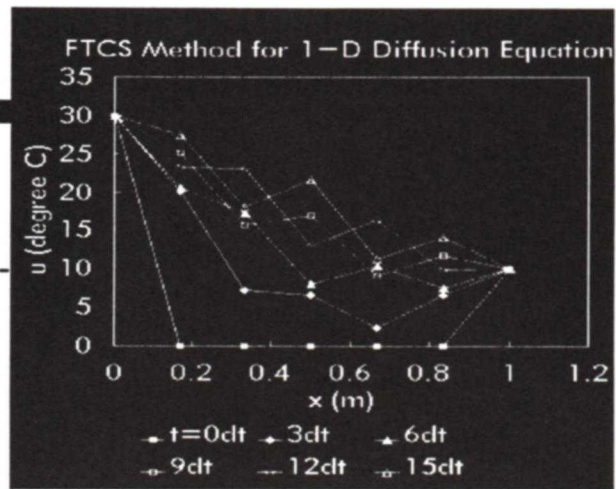
READY



(b) $F = 0.51$.

```

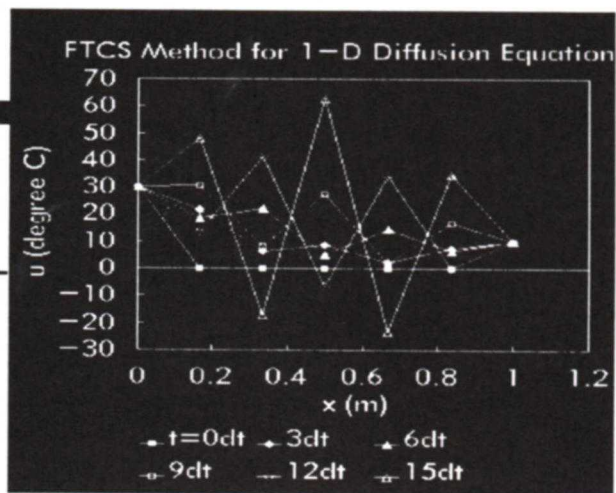
A      C      D      E
9      Constant c, c =      1.1
10     domain in x, L =      1
11     intervals, I =      6
12     Grid number, F =      0.55
13     Step in x, dx =      0.166667
14     Step in t, dt =      0.012626
15
16
17     x=      x=      x=
18     0.166667 0.333333 0.5
19 -----
20     0      0      0
21     16.5     0      0
22     14.85    9.075   0
23     20.00625 7.26    6.655
24     18.49238 13.93769 4.6585
25     22.31649 11.33921 11.09721
26     20.50492 17.24362 8.010956
27     23.9335  13.95937 14.48762
28     21.7843  19.69168 10.3583
FTCS.WK3
    
```



(c) $F = 0.55$

```

A      C      D      E
9      Constant c, c =      1.1
10     domain in x, L =      1
11     intervals, I =      6
12     Grid number, F =      0.6
13     Step in x, dx =      0.166667
14     Step in t, dt =      0.013774
15
16
17     x=      x=      x=
18     0.166667 0.333333 0.5
19 -----
20     0      0      0
21     18      0      0
22     14.4     10.8    0
23     21.6     6.48   8.64
24     17.568   16.848  3.456
25     24.5952  9.2448  14.8608
26     18.62784 21.82464 4.8384
27     27.36922 9.714816 20.72218
28     18.35505 26.91187 3.942605
FTCS.WK3
    
```



(d) $F = 0.6$

```

A      A      B      C
7      du(x,t)/dt = c^2 x d^2u(x,t)/
8
9      Constant
10     Solution domain in
11     Number of x interval
12     Grid Fourier number
13     Step in
14     Step in
15
16     RESULTS
17     x=      x=
18     t      0 0.166667
19 -----
20     0      10     0
21     0.006887 10     3
22     0.013774 10     4.2
23     0.020661 10     4.95
24     0.027548 10     5.466
25     0.034435 10     5.8749
26     0.041322 10     6.22422
FTCS.WK3
    
```

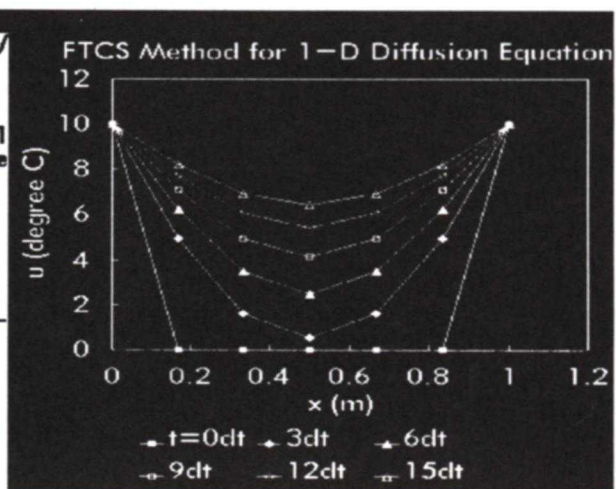


Fig. 6. The effects of changing the boundary conditions to $u(0, t) = 10$ using FTCS.WK3 ($I = 6, F = 0.3$)

A:C20: -20*@SIN(@PI*C18)

READY

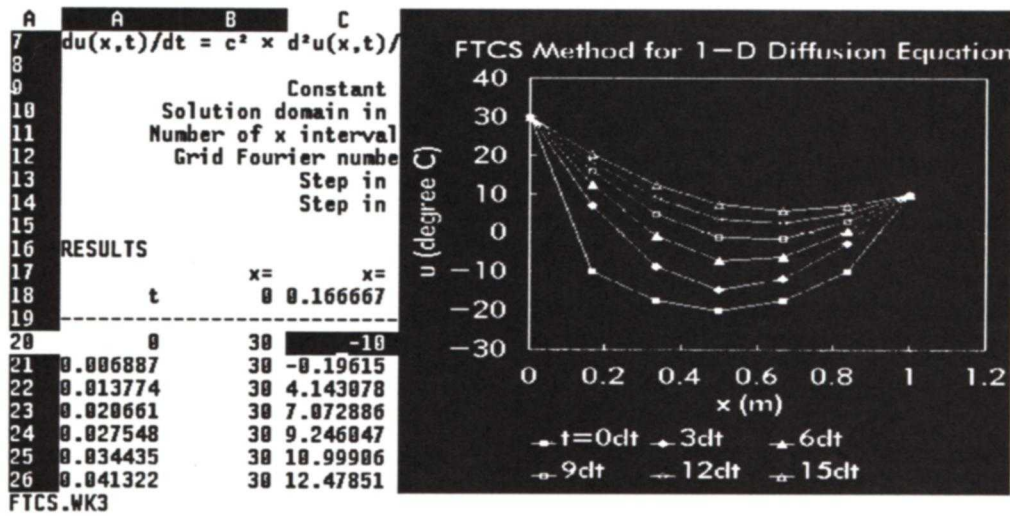


Fig. 7. The effects of changing the initial condition to $u(x, 0) = -20\sin(\pi x)$ using FTCS.WK3 ($I = 6, F = 0.3$)

A:B29: 60

READY

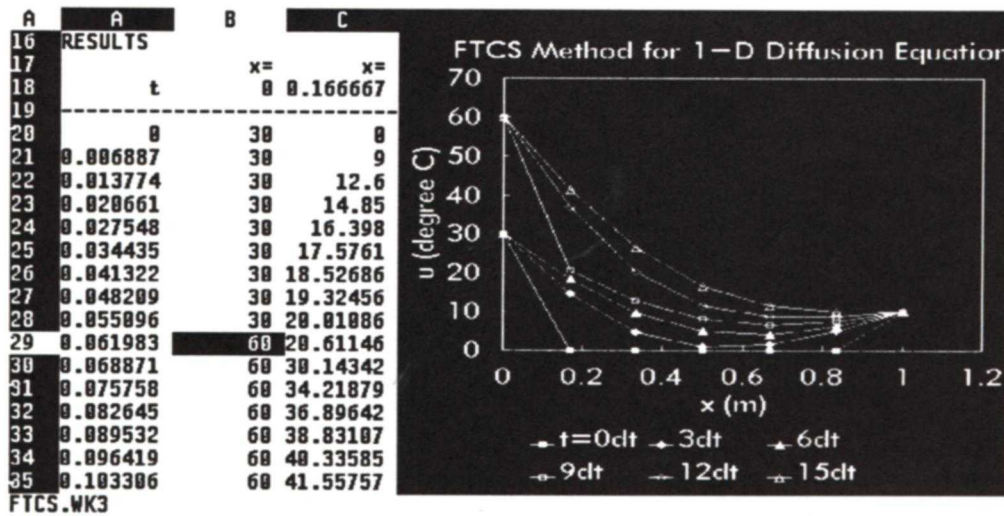


Fig. 8. The effects of changing to a variable boundary condition $u(0, t) = 30$ for $t < 9\Delta t$, $u(0, t) = 60$ for $t \geq 9\Delta t$ using FTCS.WK3 ($I = 6, F = 0.3$)

at the corresponding x -grid points using the built-in mathematical functions of 1-2-3. The alternative way is to directly enter the pre-calculated numerical values of the initial conditions at the x -grid points. By using this technique, problems with variable boundary conditions can be easily treated. Figure 8 displays the results when the original boundary condition $u(0, t)$ is suddenly increased from 30 to 60°C at $t = 9\Delta t$.

The spreadsheet program CN.WK3

This spreadsheet program employs the CN method to solve the diffusion equation. To march one t -step, equation (4) is solved by the matrix

inversion technique. This is done by the 1-2-3's /Data Matrix Invert and /Data Matrix Multiply commands programmed in the macros.

Example. Consider the same problem solved previously.

The input data are the same as those shown in Fig. 2 for FTCS.WK3. In this case, the computed results for 15 t -steps and the variations of u with x are obtained as shown in Fig. 9. For CN.WK3, the interactive graphics mode cannot be used as explained previously. However, the results for $F = 0.6$ are obtained as shown in Fig. 10 which does not exhibit any numerical instability.

The spreadsheet program CD.WK3

This spreadsheet program employs the explicit CD method to solve the wave equation. To march one t -step, the program assigns equation (6a) or (6b) accordingly to the interior grid points at the end of the t -step and the numerical results are obtained. This is done by the 1-2-3's /Copy command programmed in the macros.

Example. Consider the problem of finding the deflection $u(x, t)$ of an elastic string of length $L = 1$ m and $c = 1$ m/s. The string is held fixed at both ends and released from rest with the initial deflection $u(x, 0) = \sin(\pi x)/15$.

The input data c, L, I, R , the boundary conditions and the initial conditions are entered as shown in Fig. 11. The computed results for 15 t -steps are obtained as shown in Fig. 12(a), the variation of u with x and the variation of u with t are shown in Fig. 12(b-c). From Fig. 12(c), it is easily seen that the period of the vibration is 2 s.

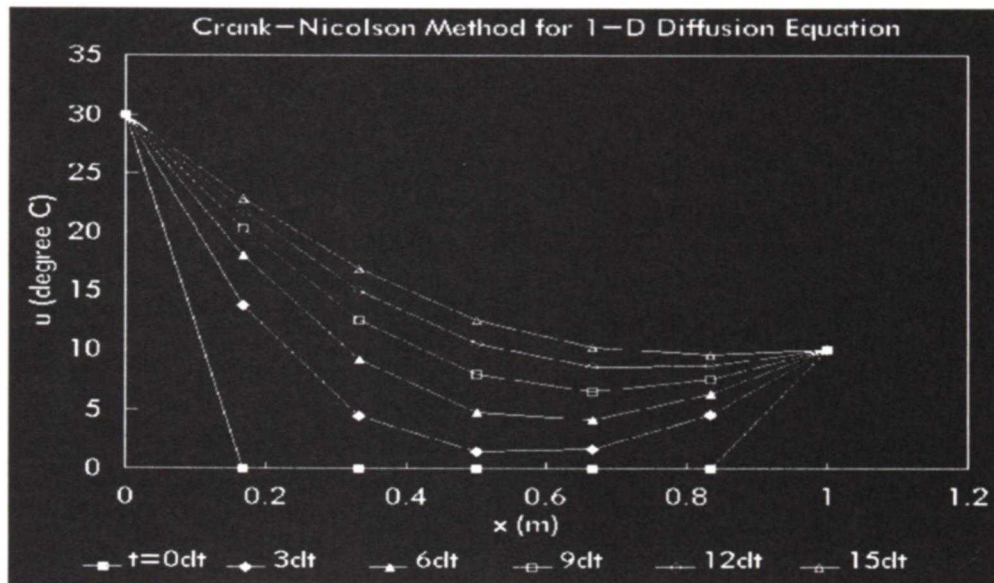
The interactive graphics mode is used to perform the 'What-if' analysis by following the same procedure as that described for FTCS.WK3. The effects of varying the Courant number R are shown in Figs 13 and 14. These figures clearly demonstrate that the CD method is unstable for $R > 1$. Figure 15 displays the results by changing the initial condition

A:A35: +A34*DT READY
 Proceed another t-step ? <Y/N> : _

| A | B | C | D | E | F | G | H | |
|----|----------|----|----------|----------|----------|----------|----------|----|
| 16 | Results | | | | | | | |
| 17 | | x= | x= | x= | x= | x= | x= | |
| 18 | t | 0 | 0.166667 | 0.333333 | 0.5 | 0.666667 | 0.833333 | 1 |
| 19 | | | | | | | | |
| 20 | 0 | 30 | 0 | 0 | 0 | 0 | 10 | |
| 21 | 0.006887 | 30 | 7.018219 | 0.824565 | 0.128808 | 0.284834 | 2.340558 | 10 |
| 22 | 0.013774 | 30 | 11.09949 | 2.619332 | 0.60719 | 0.936215 | 3.708883 | 10 |
| 23 | 0.020661 | 30 | 13.72261 | 4.512353 | 1.454215 | 1.701745 | 4.609163 | 10 |
| 24 | 0.027548 | 30 | 15.55581 | 6.265784 | 2.513184 | 2.514704 | 5.276063 | 10 |
| 25 | 0.034435 | 30 | 16.92614 | 7.833681 | 3.656446 | 3.346841 | 5.824981 | 10 |
| 26 | 0.041322 | 30 | 18.00547 | 9.225091 | 4.805544 | 4.179018 | 6.312589 | 10 |
| 27 | 0.048209 | 30 | 18.8899 | 10.46183 | 5.917891 | 4.99657 | 6.7655 | 10 |
| 28 | 0.055096 | 30 | 19.63623 | 11.56599 | 6.97268 | 5.788673 | 7.195105 | 10 |
| 29 | 0.061983 | 30 | 20.27982 | 12.5567 | 7.961357 | 6.54789 | 7.605429 | 10 |
| 30 | 0.068871 | 30 | 20.84373 | 13.44978 | 8.881953 | 7.269554 | 7.997244 | 10 |
| 31 | 0.075758 | 30 | 21.3437 | 14.25818 | 9.735896 | 7.951132 | 8.370134 | 10 |
| 32 | 0.082645 | 30 | 21.79091 | 14.9925 | 10.52627 | 8.59166 | 8.723471 | 10 |
| 33 | 0.089532 | 30 | 22.19364 | 15.66146 | 11.25686 | 9.191289 | 9.056825 | 10 |
| 34 | 0.096419 | 30 | 22.55817 | 16.27237 | 11.9317 | 9.750942 | 9.370086 | 10 |
| 35 | 0.103306 | 30 | 22.88945 | 16.83136 | 12.55477 | 10.27205 | 9.663468 | 10 |

CN.WK3 CMD

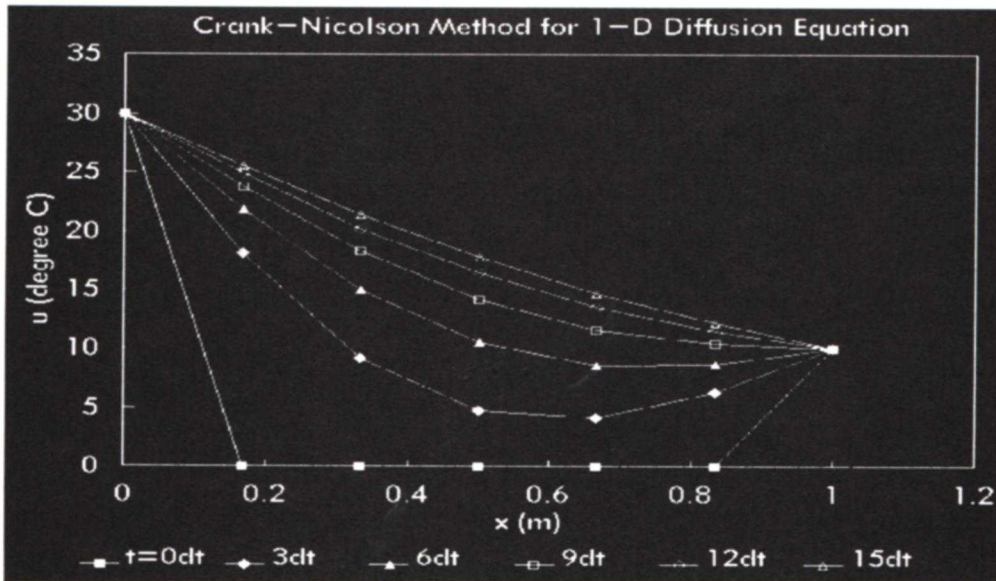
Fig. 9. The solution of the one-dimensional heat conduction problem for $I = 6$ and $F = 0.3$ using CN.WK3
 (a) The solution for 15 t -steps



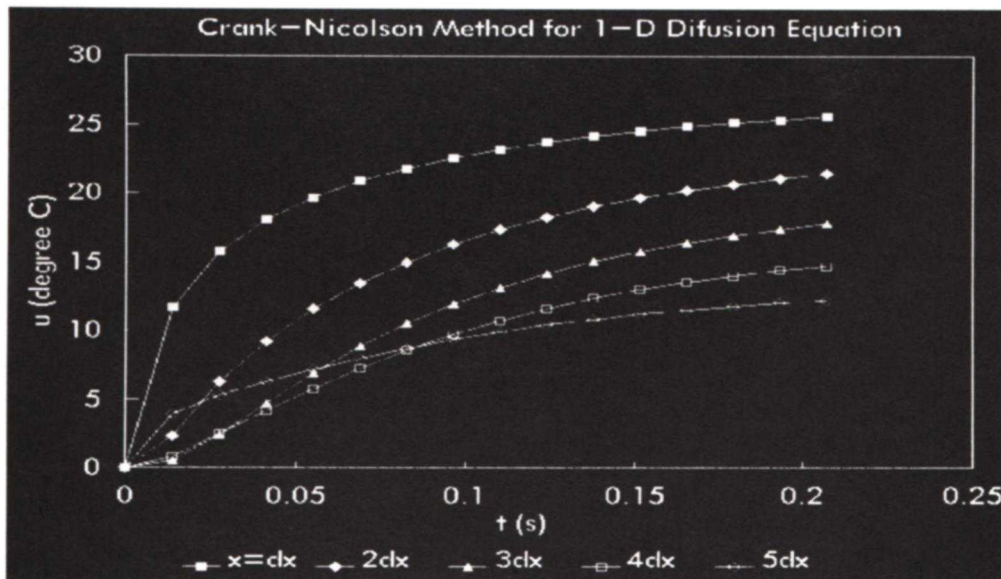
(b) Variations of u with x as displayed on the screen

| A | B | C | D | E | F | G | H | |
|----|----------|-------|----------|----------|----------|----------|----------|----|
| 16 | Results | | | | | | | |
| 17 | | x= | x= | x= | x= | x= | x= | |
| 18 | t | 0 | 0.166667 | 0.333333 | 0.5 | 0.666667 | 0.833333 | 1 |
| 19 | | ----- | | | | | | |
| 20 | 0 | 30 | 0 | 0 | 0 | 0 | 10 | |
| 21 | 0.013774 | 30 | 11.6814 | 2.380797 | 0.58952 | 0.843307 | 3.90812 | 10 |
| 22 | 0.027548 | 30 | 15.77958 | 6.281755 | 2.384465 | 2.505263 | 5.354887 | 10 |
| 23 | 0.041322 | 30 | 18.18979 | 9.264369 | 4.760458 | 4.158436 | 6.338165 | 10 |
| 24 | 0.055096 | 30 | 19.68991 | 11.60212 | 6.965303 | 5.776082 | 7.197264 | 10 |
| 25 | 0.068871 | 30 | 20.87523 | 13.47924 | 8.889326 | 7.265231 | 7.994562 | 10 |
| 26 | 0.082645 | 30 | 21.81174 | 15.01638 | 10.53876 | 8.593457 | 8.722144 | 10 |
| 27 | 0.096419 | 30 | 22.57328 | 16.29212 | 11.94569 | 9.7567 | 9.371191 | 10 |
| 28 | 0.110193 | 30 | 23.20317 | 17.36843 | 13.14399 | 10.76443 | 9.94051 | 10 |
| 29 | 0.123967 | 30 | 23.72968 | 18.26029 | 14.16415 | 11.63168 | 10.4344 | 10 |
| 30 | 0.137741 | 30 | 24.17269 | 19.02115 | 15.03256 | 12.37497 | 10.85985 | 10 |
| 31 | 0.151515 | 30 | 24.54702 | 19.66604 | 15.77175 | 13.01041 | 11.22472 | 10 |
| 32 | 0.165289 | 30 | 24.86417 | 20.21349 | 16.40094 | 13.55276 | 11.53677 | 10 |
| 33 | 0.179063 | 30 | 25.13332 | 20.67867 | 16.93651 | 14.01521 | 11.80319 | 10 |
| 34 | 0.192837 | 30 | 25.362 | 21.07421 | 17.39238 | 14.40926 | 12.03038 | 10 |
| 35 | 0.206612 | 30 | 25.55641 | 21.41065 | 17.78041 | 14.74491 | 12.224 | 10 |

Fig. 10. The solution of the one-dimensional heat conduction problem for $I = 6$ and $F = 0.6$ using CN.WK3.
 (a) The solution for 15 t -steps



(b) Variations of u with x as displayed on the screen



(c) Variations of u with t as displayed on the screen

```

A:H20: 0
Any changes for u(x,0) ? <Y/N> : _
READY
A      A      B      C      D      E      F      G      H
1 =====
2 SOLVING THE 1-D WAVE EQUATION BY THE EXPLICIT CENTRAL DIFFERENCE METHOD
3 Conditions : constant steps in x and t, constant B.C.'s (may not equal)
4 [© C Y Lam, Manyang Technological University]
5 =====
6 Equation to be solved :
7  $d^2u(x,t)/dt^2 = c^2 \times d^2u(x,t)/dx^2$  where d denotes partial differentiation
8
9      Constant c, c =          1
10     Solution domain in x, L =      1
11     Number of x intervals, I =     6
12     Courant number, R =           1
13     Step in x, dx = 0.166667
14     Step in t, dt = 0.166667
15
16 Boundary and initial conditions u(x,0) at t=0 are :
17      x=      x=      x=      x=      x=      x=
18      t      0 0.166667 0.333333      0.5 0.666667 0.833333      1
19 -----
20      0      0 0.033333 0.057735 0.066667 0.057735 0.033333      0
CD.WK3
CMD
    
```

Fig. 11. The input data in CD.WK3
 (a) The input data and the boundary and initial conditions $u(x, 0)$

```

A:H26: 0
Any changes for du(x,0)/dt ? <Y/N> : _
READY
A      A      B      C      D      E      F      G      H
22 Initial conditions du(x,0)/dt at t=0 are :
23      x=      x=      x=      x=      x=      x=      x=
24      t      0 0.166667 0.333333      0.5 0.666667 0.833333      1
25 -----
26      0      0      0      0      0      0      0      0
27
    
```

(b) The initial condition $\partial u(x, 0)/\partial t$

$u(x, 0)$ to $u(x, 0) = 0.015x$ for $0 \leq x \leq 4\Delta x$ and $u(x, 0) = 0.03(1 - x)$ for $4\Delta x \leq x \leq 1$. The other initial and boundary conditions can also be altered to study the effects of the changing.

CONCLUSIONS

Three spreadsheet programs have been developed to solve the one-dimensional diffusion equation by the explicit FTCS method and the implicit CN method, and the one-dimensional wave equation by the explicit CD method. These programs are menu-driven, user-friendly and interactive. Little spreadsheet knowledge is required to use these programs. The programs use constant Dirichlet boundary conditions. However, variable boundary conditions can also be used if they are entered manually. The powerful interactive graphics feature allows the user to carry out a series of numerical experiments easily by varying different input

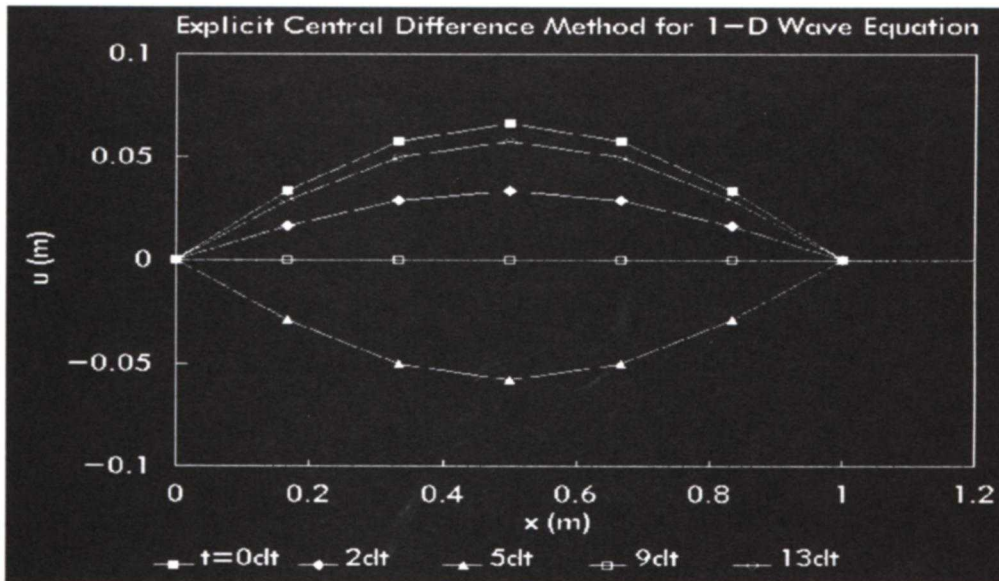
parameters and to view the effects automatically in graphical form. These features, which are difficult or impossible to obtain by the traditional programming approach, certainly help the students to improve their understanding of the numerical methods. The limitations of these spreadsheet programs are similar to those of the programs developed for Laplace equation as discussed in Part 1 [1]. They include a maximum of six curves that can be graphed together and the number of x -grid points cannot be too large to avoid memory full error. However, the limitations do not present any problems for educational purposes as the grid systems need not be very fine. Classroom use of these programs has resulted favourable comments by the students, especially in the ease of use and the interactive graphics feature. The spreadsheet approach can be extended to other finite difference methods, other problems subject to Neumann boundary conditions and other parabolic and hyperbolic equations. Work in this direction is being pursued.

A:A47: +A46+\$DT
 Proceed another t-step ? <Y/N> : _

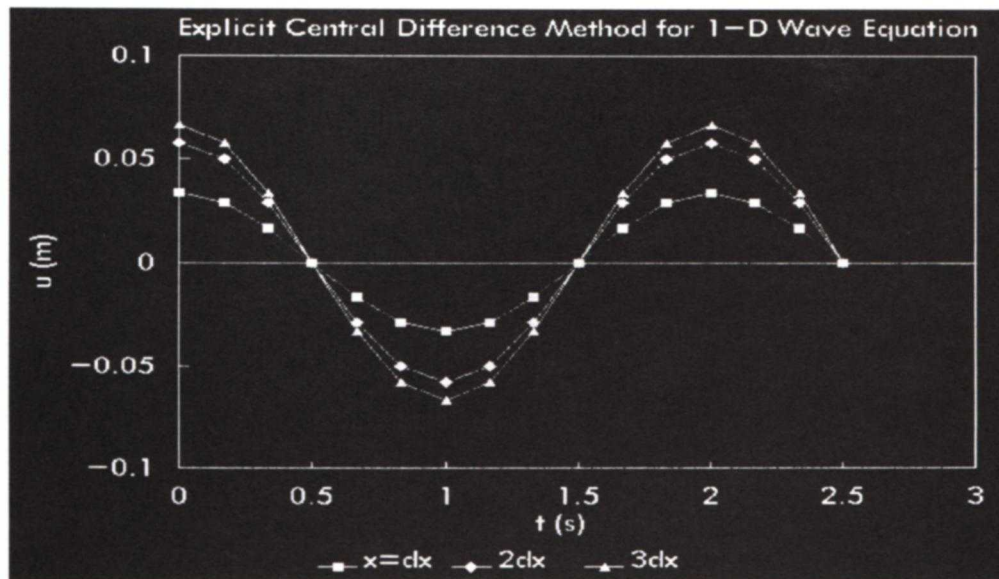
READY

| A | B | C | D | E | F | G | H | |
|----|----------|----|----------|----------|----------|----------|----------|---|
| 28 | RESULTS | | | | | | | |
| 29 | | x= | x= | x= | x= | x= | x= | |
| 30 | t | 0 | 0.166667 | 0.333333 | 0.5 | 0.666667 | 0.833333 | 1 |
| 31 | | | | | | | | |
| 32 | 0 | 0 | 0.033333 | 0.057735 | 0.066667 | 0.057735 | 0.033333 | 0 |
| 33 | 0.166667 | 0 | 0.028868 | 0.05 | 0.057735 | 0.05 | 0.028868 | 0 |
| 34 | 0.333333 | 0 | 0.016667 | 0.028868 | 0.033333 | 0.028868 | 0.016667 | 0 |
| 35 | 0.5 | 0 | -3.4E-21 | -6.8E-21 | -3.4E-21 | 1.0E-20 | 0 | 0 |
| 36 | 0.666667 | 0 | -0.01667 | -0.02887 | -0.03333 | -0.02887 | -0.01667 | 0 |
| 37 | 0.833333 | 0 | -0.02887 | -0.05 | -0.05774 | -0.05 | -0.02887 | 0 |
| 38 | 1 | 0 | -0.03333 | -0.05774 | -0.06667 | -0.05774 | -0.03333 | 0 |
| 39 | 1.166667 | 0 | -0.02887 | -0.05 | -0.05774 | -0.05 | -0.02887 | 0 |
| 40 | 1.333333 | 0 | -0.01667 | -0.02887 | -0.03333 | -0.02887 | -0.01667 | 0 |
| 41 | 1.5 | 0 | -3.4E-21 | 3.4E-21 | -3.4E-21 | 2.0E-20 | 3.4E-21 | 0 |
| 42 | 1.666667 | 0 | 0.016667 | 0.028868 | 0.033333 | 0.028868 | 0.016667 | 0 |
| 43 | 1.833333 | 0 | 0.028868 | 0.05 | 0.057735 | 0.05 | 0.028868 | 0 |
| 44 | 2 | 0 | 0.033333 | 0.057735 | 0.066667 | 0.057735 | 0.033333 | 0 |
| 45 | 2.166667 | 0 | 0.028868 | 0.05 | 0.057735 | 0.05 | 0.028868 | 0 |
| 46 | 2.333333 | 0 | 0.016667 | 0.028868 | 0.033333 | 0.028868 | 0.016667 | 0 |
| 47 | 2.5 | 0 | -1.0E-20 | -2.7E-20 | 3.4E-21 | -1.0E-20 | 6.8E-21 | 0 |

Fig. 12. The solution of the one-dimensional heat conduction problem for $I = 6$ and $R = 1$ using CD.WK3
 (a) The solution for 15 t -steps



(b) Variations of u with x as displayed on the screen



(c) Variations of u with t as displayed on the screen

```

A      C      D      E
1      =====
2      VE EQUATION BY THE EXPLICIT
3      ant steps in x and t, const
4      g Technological University]
5      =====
6      ved :
7      x d2u(x,t)/dx2 where d deno
8
9      Constant c, c =      1
10     omain in x, L =     1
11     intervals, I =      6
12     ant number, R =     1.3
13     Step in x, dx =    0.166667
14     Step in t, dt =    0.216667
15
16     al conditions u(x,0) at t=0
17     x=      x=      x=
18     0.166667 0.333333 0.5
19
20     0.0333333 0.057735 0.066667
CD.WK3
    
```

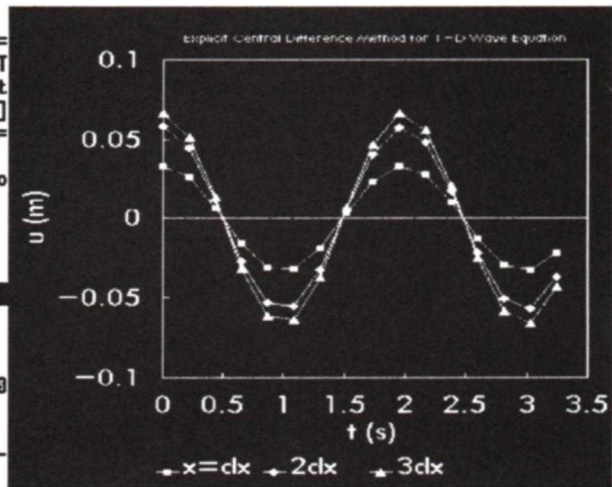
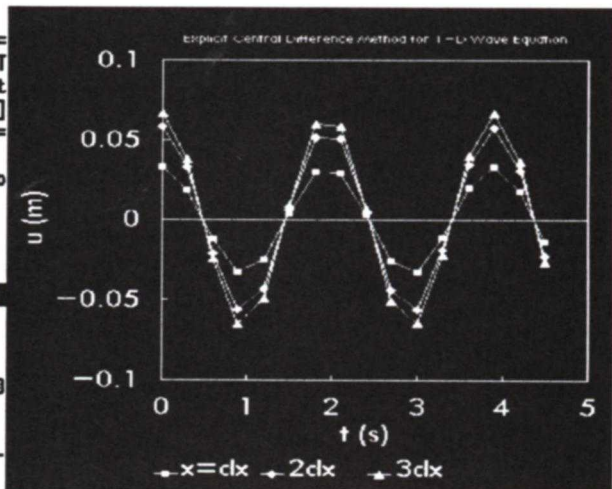


Fig. 13. The effects of varying R using CD.WK3
(a) R = 1.3

```

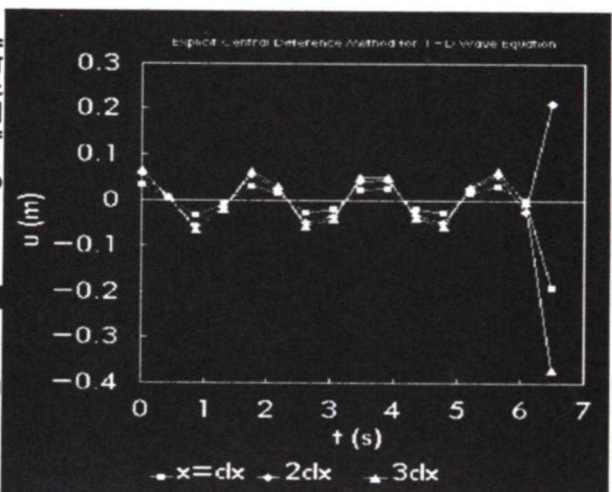
A      C      D      E
1      =====
2      VE EQUATION BY THE EXPLICIT
3      ant steps in x and t, const
4      g Technological University]
5      =====
6      ved :
7      x d2u(x,t)/dx2 where d deno
8
9      Constant c, c =      1
10     omain in x, L =     1
11     intervals, I =      6
12     ant number, R =     1.8
13     Step in x, dx =    0.166667
14     Step in t, dt =    0.3
15
16     al conditions u(x,0) at t=0
17     x=      x=      x=
18     0.166667 0.333333 0.5
19
20     0.0333333 0.057735 0.066667
CD.WK3
    
```



(b) R = 1.8

```

A      C      D      E
1      =====
2      VE EQUATION BY THE EXPLICIT
3      ant steps in x and t, const
4      g Technological University]
5      =====
6      ved :
7      x d2u(x,t)/dx2 where d deno
8
9      Constant c, c =      1
10     omain in x, L =     1
11     intervals, I =      6
12     ant number, R =     2.6
13     Step in x, dx =    0.166667
14     Step in t, dt =    0.433333
15
16     al conditions u(x,0) at t=0
17     x=      x=      x=
18     0.166667 0.333333 0.5
19
20     0.0333333 0.057735 0.066667
CD.WK3
    
```



(c) R = 2.6

A:E12: 2.9

READY

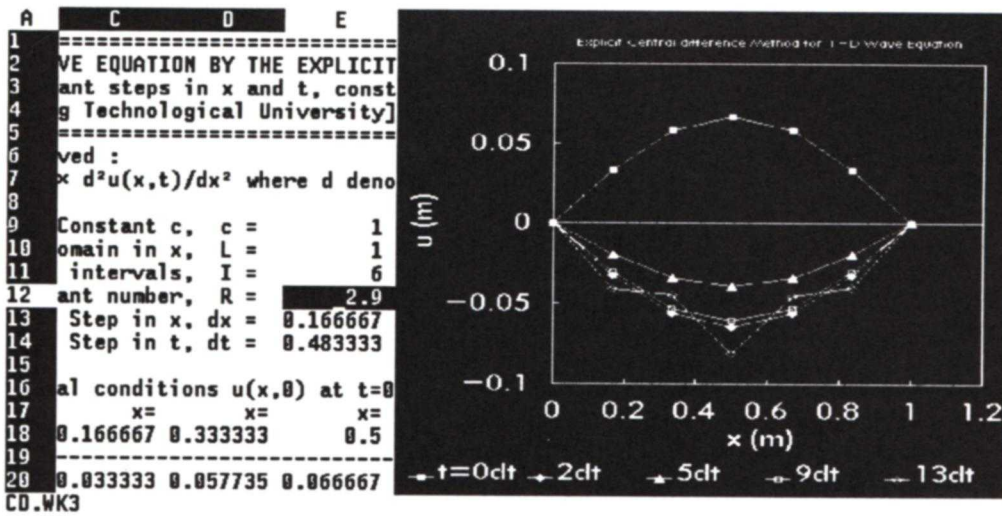


Fig. 14. The results for $R = 2.9$ using CD.WK3

A:G20: 0.03*(1-G18)

READY

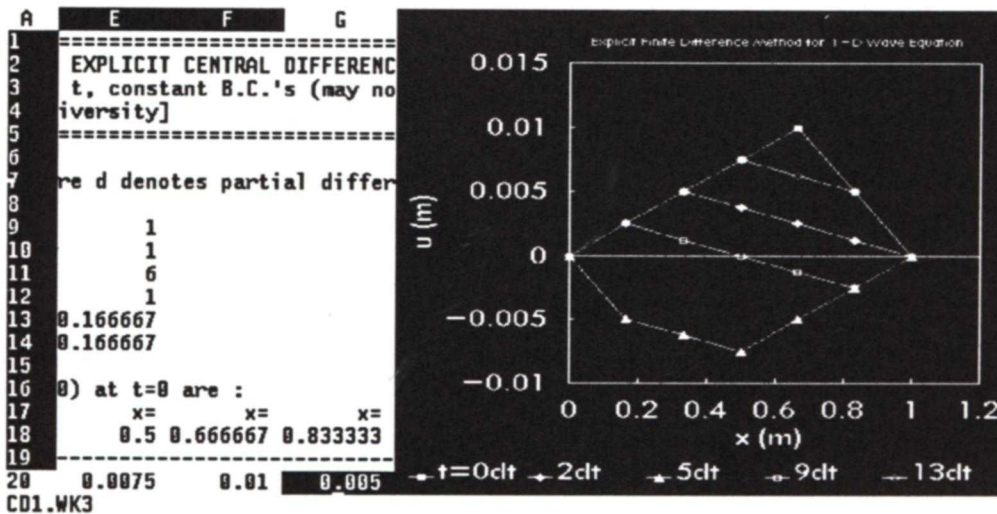


Fig. 15. The effects of changing the initial condition $u(x, 0)$ to $u(x, 0) = 0.015x$ for $0 \leq x \leq 4\Delta x$, $u(x, 0) = 0.03(1 - x)$ for $4\Delta x \leq x \leq 1$ using CD.WK3

REFERENCES

1. C. Y. Lam, Spreadsheet approach to partial differential equations. Part 1: Elliptic equation, *Int. J. Engng Educ.*, **8**, 278-287 (1992).
2. Lotus Development Corporation, *Lotus 1-2-3 Release 3.0 Reference*. Lotus Development Corporation, Cambridge (1989).
3. J. H. Ferziger, *Numerical Methods for Engineering Application*. Wiley-Interscience, New York (1981).
4. W. H. Press, B. P. Flannery, S. A. Teukolsky and W. T. Vetterling, *Numerical Recipes—The Art of Scientific Computing (Fortran Version)*. Cambridge University Press, Cambridge (1989).
5. J. D. Hoffman, *Numerical Methods for Engineers and Scientists*. McGraw-Hill, New York (1992).